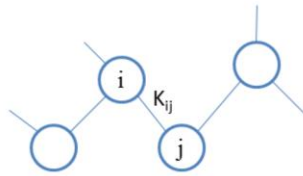


Equation of motion for coupled oscillator

$$F_{ix} = m_i \frac{d^2}{dt^2} x_i$$

$$= \sum_j \frac{d^2 \phi}{dx_i dx_j} (x_j - x_i) \equiv \sum_j k_{ij} (x_j - x_i)$$



$$\left. \begin{aligned} K_{ij} &= -k_{ij} \\ K_{ii} &= \sum_j k_{ij} \end{aligned} \right\} \Rightarrow F_{ix} = -\sum_j K_{ij} x_j$$

$$x_i = x_{0i} \exp(-i\omega t)$$

$$\Rightarrow F_{ix} = -m_i \omega^2 x_i$$

$$\sum_j K_{ij} x_j = m_i \omega^2 x_i \Rightarrow Kx = Mx\omega^2$$

$$V = \frac{1}{2} x^t Kx = \frac{1}{2} x^t Mx\omega^2$$

$\phi$ : two-body potential

When a force proportional to displacement is taken into account, it is called harmonic approximation.

**K**: stiffness matrix

Oscillating  $x$  is assumed as a solution.

Simultaneous equations for all the particles.

→ An equation of matrix

$V$ : Vibrational energy

Hessian method

$$M^{-1}Kx = x\omega^2$$

$$M = \begin{pmatrix} m_1 & & & 0 \\ & m_1 & & \\ & & \ddots & \\ 0 & & & m_N \\ & & & & m_N \\ & & & & & m_N \end{pmatrix}$$

$$K = \begin{pmatrix} \frac{\partial^2 U}{\partial \alpha_1^2} & \frac{\partial^2 U}{\partial \alpha_1 \partial \beta_1} & \frac{\partial^2 U}{\partial \alpha_1 \partial \gamma_1} & & & & \\ \frac{\partial^2 U}{\partial \beta_1 \partial \alpha_1} & \frac{\partial^2 U}{\partial \beta_1^2} & \frac{\partial^2 U}{\partial \beta_1 \partial \gamma_1} & \dots & & & \\ \frac{\partial^2 U}{\partial \gamma_1 \partial \alpha_1} & \frac{\partial^2 U}{\partial \gamma_1 \partial \beta_1} & \frac{\partial^2 U}{\partial \gamma_1^2} & & \dots & & \\ \vdots & & & \ddots & & & \\ & & & & \ddots & & \\ & & & & & \frac{\partial^2 U}{\partial \alpha_N^2} & \frac{\partial^2 U}{\partial \alpha_N \partial \beta_N} & \frac{\partial^2 U}{\partial \alpha_N \partial \gamma_N} \\ \vdots & & & & & \frac{\partial^2 U}{\partial \beta_N \partial \alpha_N} & \frac{\partial^2 U}{\partial \beta_N^2} & \frac{\partial^2 U}{\partial \beta_N \partial \gamma_N} \\ & & & & & \frac{\partial^2 U}{\partial \gamma_N \partial \alpha_N} & \frac{\partial^2 U}{\partial \gamma_N \partial \beta_N} & \frac{\partial^2 U}{\partial \gamma_N^2} \end{pmatrix}$$

$$M^{-1} = \begin{pmatrix} \mu_1 & & & 0 \\ & \mu_1 & & \\ & & \ddots & \\ 0 & & & \mu_N \\ & & & & \mu_N \\ & & & & & \mu_N \end{pmatrix}, x = \begin{pmatrix} x_1 \\ y_1 \\ z_1 \\ \vdots \\ x_N \\ y_N \\ z_N \end{pmatrix}$$

$x$  is an eigenvector of  $M^{-1}K$ .

→ You have to find a set of  $x$  that diagonalizes  $M^{-1}K$ .

### GF method

A variation of Hessian method, where the coordinate system is transformed so that the variables represents molecular parameters (such as bond length, bond angle, etc.).

$$\begin{cases} M^{-1}Kx = x\omega^2 \\ x' = U^t x \end{cases}$$

$$\Rightarrow U^t M^{-1} U U^t K U x' = U^t U x' \omega^2$$

$$\begin{cases} U^t M^{-1} U \equiv G \\ U^t K U \equiv F \end{cases}$$

$$\Rightarrow G F x' = x' \omega^2$$

U transforms the Cartesian coordinate into internal molecular coordinate.

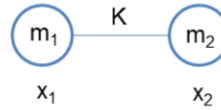
F is a set of parameters such as stretching, biting, etc., intuitively understandable.

$x'$  is an eigenvector of GF.

→ You have to find a set of  $x'$  that diagonalizes GF matrix.

Natural oscillation of a diatomic molecule

$$\begin{pmatrix} \mu_1 & 0 \\ 0 & \mu_2 \end{pmatrix} \begin{pmatrix} K & -K \\ -K & K \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} \omega^2$$



Hesseian Method

$$\begin{pmatrix} \mu_1 K - \omega^2 & -\mu_1 K \\ -\mu_2 K & \mu_2 K - \omega^2 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = 0$$

when  $\omega^2 = 0$

$$x_1 = \frac{1}{\sqrt{2}}, \quad x_2 = \frac{1}{\sqrt{2}}$$

when  $\omega^2 = (\mu_1 + \mu_2)K$

$$x_1 = \frac{\mu_1}{\sqrt{\mu_1^2 + \mu_2^2}}, \quad x_2 = -\frac{\mu_2}{\sqrt{\mu_1^2 + \mu_2^2}}$$

GF method

$$x' = U^T x = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} x_1 + x_2 \\ x_1 - x_2 \end{pmatrix}$$

$$G = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} \mu_1 & 0 \\ 0 & \mu_2 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} \mu_1 + \mu_2 & \mu_1 - \mu_2 \\ \mu_1 - \mu_2 & \mu_1 + \mu_2 \end{pmatrix}$$

$$F = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} K & -K \\ -K & K \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 2K \end{pmatrix}$$

$$GF = \begin{pmatrix} 0 & (\mu_1 - \mu_2)K \\ 0 & (\mu_1 + \mu_2)K \end{pmatrix}$$

$$\begin{pmatrix} -\omega^2 & (\mu_1 - \mu_2)K \\ 0 & (\mu_1 + \mu_2)K - \omega^2 \end{pmatrix} \begin{pmatrix} x'_1 \\ x'_2 \end{pmatrix} = 0$$

### Mass-weighted Hessian method

$$M^{-1}Kx = x\omega^2$$

$$M^{-1/2}KM^{-1/2}M^{1/2}x = M^{1/2}x\omega^2$$

$$M^{1/2} = \begin{pmatrix} \sqrt{m_1} & & & & & & & 0 \\ & \sqrt{m_1} & & & & & & \\ & & \sqrt{m_1} & & & & & \\ & & & \dots & & & & \\ & & & & \sqrt{m_N} & & & \\ & 0 & & & & \sqrt{m_N} & & \\ & & & & & & \sqrt{m_N} & \\ & & & & & & & \sqrt{m_N} \end{pmatrix}$$

$M^{-1}K$  is not a symmetric matrix.  
→  $x$ 's are not orthogonal to each other.  
→  $\omega$  is not necessarily real.

When we multiply  $M^{1/2}$  on the both hands of the equation.

$M$  is a diagonal matrix.  
→  $M^{1/2}$  is easily obtained.

$$\left. \begin{array}{l} M^{-1/2}KM^{-1/2} \equiv D \\ M^{1/2}x \equiv w \end{array} \right\} \Rightarrow Dw = w\omega^2 \quad D_{ij} = \sqrt{\mu_i}\sqrt{\mu_j}K_{ij}$$

$D$  is symmetrical matrix.  
→ so-called dynamical matrix

Natural oscillation of a diatomic molecule

$$\begin{pmatrix} \mu_1 K & -\sqrt{\mu_1 \mu_2} K \\ -\sqrt{\mu_1 \mu_2} K & \mu_2 K \end{pmatrix} \begin{pmatrix} w_1 \\ w_2 \end{pmatrix} = \begin{pmatrix} w_1 \\ w_2 \end{pmatrix} \omega^2$$

$$\begin{pmatrix} \mu_1 K - \omega^2 & -\sqrt{\mu_1 \mu_2} K \\ -\sqrt{\mu_1 \mu_2} K & \mu_2 K - \omega^2 \end{pmatrix} \begin{pmatrix} w_1 \\ w_2 \end{pmatrix} = 0$$

when  $\omega^2 = 0$

$$w_1 = \sqrt{\frac{\mu_2}{\mu_1 + \mu_2}}, \quad w_2 = \sqrt{\frac{\mu_1}{\mu_1 + \mu_2}}$$

when  $\omega^2 = (\mu_1 + \mu_2)K$

$$w_1 = \sqrt{\frac{\mu_1}{\mu_1 + \mu_2}}, \quad w_2 = -\sqrt{\frac{\mu_2}{\mu_1 + \mu_2}}$$

Compared with GF method

$$GFx' = x'\omega^2$$

$$G^{1/2}FG^{1/2}G^{-1/2}x'$$

$$\left\{ \begin{array}{l} G^{1/2}FG^{1/2} \equiv D' \\ G^{-1/2}x' \equiv w' \end{array} \right.$$

$$\Rightarrow D'w' = w'\omega^2$$

$$\left\{ \begin{array}{l} w' = G^{-1/2}x' \\ = U^t M^{1/2} U U^t M^{-1/2} w = U^t w \end{array} \right.$$

$$\left\{ \begin{array}{l} D' = G^{1/2}FG^{1/2} \\ = U^t M^{-1/2} U U^t K U U^t M^{-1/2} U \\ = U^t D U \end{array} \right.$$

Symmetrize GF matrix

→  $D$  and  $D'$  are transformed to each other using  $U$ .

Natural oscillation of a diatomic molecule (cont.)

mass-weighted coordinate (orthonormalized)

$$\hat{\mathbf{W}} = \frac{1}{\sqrt{\mu_1 + \mu_2}} \begin{pmatrix} \sqrt{\mu_2} & \sqrt{\mu_1} \\ \sqrt{\mu_1} & -\sqrt{\mu_2} \end{pmatrix}$$

Cartesian coordinate

$$\mathbf{X} = \mathbf{M}^{-1/2} \hat{\mathbf{W}} = \frac{1}{\sqrt{\mu_1 + \mu_2}} \begin{pmatrix} \sqrt{\mu_1 \mu_2} & \mu_1 \\ \sqrt{\mu_1 \mu_2} & -\mu_2 \end{pmatrix}$$

Cartesian coordinate (normalized)

$$\hat{\mathbf{X}} = \frac{1}{\sqrt{2(\mu_1^2 + \mu_2^2)}} \begin{pmatrix} \sqrt{\mu_1^2 + \mu_2^2} & \sqrt{2}\mu_1 \\ \sqrt{\mu_1^2 + \mu_2^2} & -\sqrt{2}\mu_2 \end{pmatrix}$$

modal mass matrix

$$\hat{\mathbf{X}}^T \mathbf{M} \hat{\mathbf{X}} = \begin{pmatrix} \mu_1 + \mu_2 & 0 \\ 2\mu_1 \mu_2 & \mu_1 + \mu_2 \\ 0 & \mu_1^2 + \mu_2^2 \end{pmatrix} \equiv \begin{pmatrix} m_T & 0 \\ 0 & m_V \end{pmatrix}$$

modal stiffness matrix

$$\hat{\mathbf{X}}^T \mathbf{K} \hat{\mathbf{X}} = \begin{pmatrix} 0 & 0 \\ 0 & \frac{(\mu_1 + \mu_2)^2}{\mu_1^2 + \mu_2^2} K \end{pmatrix} \equiv \begin{pmatrix} K_T & 0 \\ 0 & K_V \end{pmatrix}$$

angular frequency

$$\omega_T = \sqrt{\frac{K_T}{m_T}}, \quad \omega_V = \sqrt{\frac{K_V}{m_V}}$$

cf.

$$\omega_0 = 0, \quad \omega_1 = \sqrt{\frac{K}{m^*}} \quad \left( m^* = \frac{1}{\mu_1 + \mu_2} \right)$$

Natural oscillation of a linear triatomic molecule

$$\begin{array}{c}
 \textcircled{1} \quad K_{12} \quad \textcircled{2} \quad K_{23} \quad \textcircled{3} \\
 m_1 \quad m_2 \quad m_3
 \end{array}
 \quad \mathbf{M}^{-1/2} \mathbf{K} \mathbf{M}^{-1/2} = \begin{pmatrix} \mu_1 K_{12} & -\sqrt{\mu_1 \mu_2} K_{12} & 0 \\ -\sqrt{\mu_1 \mu_2} K_{12} & \mu_2 (K_{12} + K_{23}) & -\sqrt{\mu_2 \mu_3} K_{23} \\ 0 & -\sqrt{\mu_2 \mu_3} K_{23} & \mu_3 K_{23} \end{pmatrix}$$

$$\omega^2 = 0, \frac{1}{2} \left\{ (\mu_1 + \mu_2) K_{12} + (\mu_2 + \mu_3) K_{23} \pm \sqrt{[(\mu_1 + \mu_2) K_{12} - (\mu_2 + \mu_3) K_{23}]^2 + 4 \mu_2^2 K_{12} K_{23}} \right\}$$

$$\begin{array}{l}
 m_1 = m_3 \\
 K_{12} = K_{23} = K
 \end{array}
 \quad \rightarrow \quad \hat{\mathbf{X}} = \begin{pmatrix} \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} & \frac{\mu_1}{\sqrt{2(\mu_1^2 + 2\mu_2^2)}} \\ \frac{1}{\sqrt{3}} & 0 & -\frac{2\mu_2}{\sqrt{2(\mu_1^2 + 2\mu_2^2)}} \\ \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} & \frac{\mu_1}{\sqrt{2(\mu_1^2 + 2\mu_2^2)}} \end{pmatrix}$$

$$\hat{\mathbf{X}}^T \mathbf{M} \hat{\mathbf{X}} = \begin{pmatrix} \frac{2m_1 + m_2}{3} & 0 & 0 \\ 0 & m_1 & 0 \\ 0 & 0 & \frac{m_1 m_2 (2m_1 + m_2)}{2m_1^2 + m_2^2} \end{pmatrix}
 \quad \hat{\mathbf{X}}^T \mathbf{K} \hat{\mathbf{X}} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & K & 0 \\ 0 & 0 & \frac{(\mu_1 + 2\mu_2)}{\mu_1^2 + 2\mu_2^2} K \end{pmatrix}$$